ON THE NUMBER OF EIGENVALUES OF THE DIRAC OPERATOR IN A BOUNDED INTERVAL

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ABSTRACT. Let H_0 be the free Dirac operator and $V \geqslant 0$ be a positive potential. We study the discrete spectrum of $H(\alpha) = H_0 - \alpha V$ in the interval $(-1, 1)$ for large values of the coupling constant $\alpha > 0$. In particular, we obtain an asymptotic formula for the number of eigenvalues of $H(\alpha)$ situated in a bounded interval $[\lambda, \mu)$ as $\alpha \to \infty$.

1. STATEMENT OF THE MAIN THEOREM

Let H_0 be the free Dirac operator

$$H_0 = -i \sum_1^3 \gamma_j \frac{\partial}{\partial x_j} + \gamma_0,$$

where γ_j are 4×4 selfadjoint matrices obeying the conditions

$$\gamma_j \gamma_k + \gamma_k \gamma_j = \begin{cases} 0, & \text{if} \quad j \neq k; \\ 2\,\mathbb{I}, & \text{if} \quad j = k. \end{cases}$$

The operator H_0 is selfadjoint in the space $L^2(\mathbb{R}^3; \mathbb{C}^4)$, consisting of functions on $\mathbb{R}^3$ that take values in $\mathbb{C}^4$. The spectrum of H_0 is the set $\sigma(H_0) = (-\infty, -1] \cup [1, \infty)$.

Let $V \geqslant 0$ be a bounded potential on $\mathbb{R}^3$. Define $H(\alpha)$ to be the operator

$$H(\alpha) = H_0 - \alpha V, \qquad \alpha > 0.$$

We will always assume that $V \in L^3(\mathbb{R}^3)$. In this case, besides having a continuous spectrum that coincides with $\sigma(H_0)$, the operator $H(\alpha)$ might have a discrete spectrum in the interval $(-1, 1)$. Choose λ and μ so that $-1 < \lambda < \mu < 1$. We define $N(\alpha)$ to be the number of eigenvalues of $H(\alpha)$ inside $[\lambda, \mu)$.

Our main result is the theorem below which establishes the rate of growth of $N(\alpha)$ at infinity. The symbol f_+ denotes the positive part $f_+ = (|f| + f)/2$ of f, which can be either a real number or a real-valued function.

Theorem 1.1. *Let $V \geqslant 0$ be a bounded real-valued potential such that*

$$V(x) = \frac{\Phi(\theta)}{|x|^\nu}\Big(1 + o(1)\Big), \qquad \textit{as} \qquad |x| \to \infty,$$

where Φ is a continuous function on the unit sphere and $1 < \nu < 4/3$. Let $-1 < \lambda < \mu < 1$. Then for any $q \in (9/4, 3/\nu)$,

$$\lim_{\alpha \to \infty} \alpha^{-3/\nu + q} \int_0^\alpha N(t) t^{-q-1} dt = \frac{\nu}{3\pi^2(3 - \nu q)} \int_{\mathbb{R}^3} \Big[\Big((\Phi(\theta)|x|^{-\nu} + \mu)_+^2 - 1\Big)_+^{3/2} - \Big((\Phi(\theta)|x|^{-\nu} + \lambda)_+^2 - 1\Big)_+^{3/2}\Big] dx. \tag{1.1}$$

Remark. If $N(\alpha) \sim C\alpha^{3/\nu}$ as $\alpha \to \infty$, then the right hand side of (1.1) becomes $\dfrac{\nu}{3-\nu q} \cdot C$. Thus, the formula (1.1) determines the value of the contant C.

The question about the number of eigenvalues of the Dirac operator in a bounded interval is considered here for the first time. This theorem is new.

Perturbations $V \in L^3(\mathbb{R}^3)$ were studied in [16] by M. Klaus and, later, in [4] by M. Birman and A. Laptev. However, the object of the study in [16] and [4] was completely different from $N(\alpha)$, considered in this article. The main results of [16] and [4] imply that if $V \in L^3(\mathbb{R}^3)$, then the number $\mathcal{N}(\lambda, \alpha)$ of eigenvalues of $H(t)$ passing a regular point $\lambda \in (-1, 1)$ as t increases from 0 to α satisfies

$$\mathcal{N}(\lambda, \alpha) \sim \frac{1}{3\pi^2}\alpha^3 \int_{\mathbb{R}^3} V^3 dx, \qquad \text{as} \quad \alpha \to \infty. \tag{1.2}$$

In addition, M. Klaus proved in [16] that if $V \in L^3 \cap L^{3/2}$, then the asymptotic formula (1.2) holds even for $\lambda = 1$. In this case, $\mathcal{N}(\lambda, \alpha)$ is interpreted as the number of eigenvalues of $H(t)$ that appear at the right edge of the gap as t increases from 0 to α.

The crux of the problem. Observe that $N(\alpha) = \mathcal{N}(\mu, \alpha) - \mathcal{N}(\lambda, \alpha)$. However, since the expression on the right hand side of (1.2) does not depend on λ, this formula only implies that

$$N(\alpha) = o(\alpha^3), \qquad \text{as} \quad \alpha \to \infty.$$

In order to obtain an asymptotic formula for $N(\alpha)$ one would need to know the second term in the asymptotics of $\mathcal{N}(\lambda, \alpha)$. The second term in (1.2) has never been obtained. This explains why the problem is challenging. Another reason why the problem is challanging is that the Dirichlet-Neumann bracketing that is often used for Schrödinger operators cannot be applied to Dirac operators. To prove Theorem 1.1, one needs to develop a new machinery rich in tools that allow us to obtain the estimate of $N(\alpha)$ stated below.

Theorem 1.2. *Let $9/4 < q \leqslant 3$ and let $N(\alpha)$ be the number of eigenvalues of $H(\alpha)$ in the interval $[\lambda, \mu)$. Then*

$$\int_0^\infty N(\alpha)\alpha^{-q-1} d\alpha \leqslant C \int_{\mathbb{R}^3} V^q(x) dx$$

with a constant $C > 0$ independent of V.

Theorems 1.1 and 1.2 involve averaging of the function $N(\alpha)$. Averaging of eigenvalue counting functions also appeared in the papers [27] and [28]. However, the operators that were studied in these two papers are Schrödinger operators. These are the publications in which one discusses a periodic Schrödinger operator perturbed by a decaying potential αV. The same elliptic model is discussed in [25] , [26], and [29], but the asymptotics of $N(\alpha)$ is established in [25] , [26], and [29] without any averaging. To obtain such strong results, one has to impose very restrictive conditions on the derivatives of V. The remaining papers [1]-[3] [5], [6] , [9], [11] -[15] , [20], [24], that are devoted to Schrödinger operators, do not even deal with $N(\alpha)$. Instead of that, they deal with the number $\mathcal{N}(\lambda, \alpha)$ of eigenvalues passing the point λ.

Finally, we would like to mention the paper [10] which is related to the spectral theory of Dirac operators. However, the problems discussed in [10] are very different from the questions studied here.

2. Compact operators

For a compact operator T, the symbols $s_k(T)$ denote the singular values of T enumerated in the non-increasing order ($k \in \mathbb{N}$) and counted in accordance with their multiplicity. Observe that $s_k^2(T)$ are eigenvalues of T^*T. We set

$$n(s,T) = \#\{k: \quad s_k(T) > s\}, \qquad s > 0.$$

For a self-adjoint compact operator T we also set

$$n_\pm(s,T) = \#\{k: \quad \pm\lambda_k(T) > s\}, \qquad s > 0.$$

where $\lambda_k(T)$ are eigenvalues of T. Observe that

$$n_\pm(s_1+s_2, T_1+T_2) \leqslant n_\pm(s_1,T_1) + n_\pm(s_2,T_2), \qquad s_1, s_2 > 0.$$

A similar inequality holds for the function n. Also,

$$n(s_1 s_2, T_1 T_2) \leqslant n(s_1,T_1) + n(s_2,T_2), \qquad s_1, s_2 > 0.$$

Theorem 2.1. *Let A and B be two compact operators on the same Hilbert space. Then for any $r \in \mathbb{N}$,*

$$\sum_1^r s_k^p(A+B) \leqslant \sum_1^r s_k^p(A) + \sum_1^r s_k^p(B), \qquad \forall p \in (0,1], \tag{2.1}$$

and

$$\sum_1^r s_k^p(AB) \leqslant \sum_1^r s_k^p(A) s_k^p(B), \qquad \forall p > 0. \tag{2.2}$$

The first inequality was discovered by S. Rotfeld [22]. The second estimate is called Horn's inequality (see Section 11.6 of the book [7]).

Below we use the following notation for the positive and negative part of a self-adjoint operator T:

$$T_\pm = \frac{1}{2}(|T| \pm T).$$

Theorem 2.2. *Let $0 < p \leqslant 1$. Let $q \geqslant p$. Let A and B be two compact selfadjoint operators. Then for any $s > 0$,*

$$q\int_s^\infty \Big(n_+(t,A) - n_+(t,B)\Big) t^{q-1} dt \leqslant \|B\|^q + \sum_{k=1}^{n_+(s,A)+1} s_k^p\Big(|A|^{q/p}\mathrm{sgn}(A) - |B|^{q/p}\mathrm{sgn}(B)\Big). \tag{2.3}$$

Moreover, if $B \leqslant A$, then

$$q\int_s^\infty \Big(n_+(t,A) - n_+(t,B)\Big) t^{q-1} dt \leqslant \sum_{k=1}^{n_+(s,A)+1} s_k^p\Big(|A|^{q/p}\mathrm{sgn}(A) - |B|^{q/p}\mathrm{sgn}(B)\Big), \qquad \forall s > 0.$$

A proof of Theorem 2.2 can be found in [28].

Let H_0 and $V \geqslant 0$ be two selfadjoint operators acting on the same Hilbert space. Assume that V is bounded. For $\lambda \in \mathbb{R} \setminus \sigma(H_0)$, define the operator X_λ by

$$X_\lambda = W(H_0 - \lambda)^{-1} W, \qquad W = \sqrt{V}. \tag{2.4}$$

Two points λ and μ are said to be in the same spectral gap of H_0 provided $[\lambda, \mu] \subset \mathbb{R} \setminus \sigma(H_0)$.

Proposition 2.3. *Let $0 < p \leqslant 1$. Let $q \geqslant p$. Suppose the operators X_λ, X_μ are compact for the two points $\lambda < \mu$ that belong to the same spectral gap of H_0. Then for any $s > 0$,*

$$q \int_s^\infty \Big(n_+(t, X_\mu) - n_+(t, X_\lambda) \Big) t^{q-1} dt \leqslant \sum_{k=1}^{n_+(s, X_\mu)+1} s_k^p \Big(|X_\mu|^{q/p} \mathrm{sgn}(X_\mu) - |X_\lambda|^{q/p} \mathrm{sgn}(X_\lambda) \Big).$$

Proof. Here one needs to use the fact that $X_\lambda \leqslant X_\mu$. □

Let $\mathfrak{S}_\infty$ be the class of compact operators. Note that the condition

$$W|H_0 - \lambda_0|^{-1/2} \in \mathfrak{S}_\infty, \qquad \text{for some} \quad \lambda_0 \notin \sigma(H_0), \tag{2.5}$$

implies that operators (2.4) are compact for all $\lambda \in \mathbb{R} \backslash \sigma(H_0)$. Moreover, (2.5) implies that, for each $\alpha > 0$, the spectrum of $H(\alpha) = H_0 - \alpha V$ is discrete outside of $\sigma(H_0)$ because the difference of resolvent operators $(H(\alpha) - z)^{-1}$ and $(H_0 - z)^{-1}$ is compact for Im $z > 0$.

The following proposition is called the Birman-Schwinger principle:

Proposition 2.4. *Let H_0 and $V \geqslant 0$ be self-adjoint operators in a Hilbert space. Assume that* (2.5) *holds for some λ_0. Let $\mathcal{N}(\lambda, \alpha)$ be the number of eigenvalues of $H(t) = H_0 - tV$ passing through a regular point $\lambda \notin \sigma(H_0)$ as t increases from 0 to α. Then*

$$\mathcal{N}(\lambda, \alpha) = n_+(s, X_\lambda), \qquad \textit{for} \quad s\alpha = 1, \textit{ and } W = \sqrt{V}. \tag{2.6}$$

The idea of the proof of (2.6) is the following. First, one shows that $\lambda \in \sigma(H(\alpha))$, if and only if $\alpha^{-1} \in \sigma(W(H - \lambda)^{-1}W)$. This relation holds with multiplicities taken into account. After that, one simply uses the definition of the distribution function $n_+(s, X_\lambda)$.

Corollary 2.5. *Let H_0 and $V \geqslant 0$ be self-adjoint operators in a Hilbert space. Assume that* (2.5) *holds for some λ_0. Let $N(\alpha)$ be the number of eigenvalues of the operator $H(\alpha)$ in $[\lambda, \mu)$ contained in a gap of the spectrum $\sigma(H_0)$. Then*

$$N(\alpha) = n_+(s, X_\mu) - n_+(s, X_\lambda), \qquad s\alpha = 1. \tag{2.7}$$

Let $p > 0$. The class of compact operators T whose singular values satisfy

$$\|T\|_{\mathfrak{S}_p}^p := \sum_k s_k^p(T) < \infty$$

is called the Schatten class $\mathfrak{S}_p$.

The following statement provides a Hölder type inequality for products of compact operators that belong to different Schatten classes.

Proposition 2.6. *Let $T_1 \in \mathfrak{S}_p$ and $T_2 \in \mathfrak{S}_q$ where $p > 0$ and $q > 0$. Then $T_1T_2 \in \mathfrak{S}_r$, where $1/r = 1/p + 1/q$, and*

$$\|T_1T_2\|_{\mathfrak{S}_r} \leqslant \|T_1\|_{\mathfrak{S}_p} \|T_2\|_{\mathfrak{S}_q}.$$

A proof of this proposition can be found in [7].

Consider the following important example of an integral operator on $L^2(\mathbb{R}^d)$:

$$\big(Y u\big)(x) = (2\pi)^{-d/2} \int_{\mathbb{R}^d} a(x) e^{i\xi x} b(\xi) u(\xi) d\xi. \tag{2.8}$$

If F is the Fourier transform operator, $[a]$ and $[b]$ are operators of multiplication by the functions a and b, then

$$Y = [a]F^*[b].$$

The symbol $\mathbb{Q}$ below is used to denote the unit cube $[0, 1)^d$.

Theorem 2.7. *If a and b belong to $L^p(\mathbb{R}^d)$ with $2 \leqslant p < \infty$, then $Y \in \mathfrak{S}_p$ and*

$$\|Y\|_{\mathfrak{S}_p} \leqslant C\|a\|_{L^p}\|b\|_{L^p}.$$

If $0 < p < 2$ and

$$\sum_{n\in\mathbb{Z}^d}\big(\|a\|^p_{L^\infty(\mathbb{Q}+n)} + \|b\|^p_{L^\infty(\mathbb{Q}+n)}\big) < \infty,$$

then $Y \in \mathfrak{S}_p$ and

$$\|Y\|_{\mathfrak{S}_p} \leqslant C\Big(\sum_{n\in\mathbb{Z}^d}\|a\|^p_{L^\infty(\mathbb{Q}+n)}\Big)^{1/p}\Big(\sum_{n\in\mathbb{Z}^d}\|b\|^p_{L^\infty(\mathbb{Q}+n)}\Big)^{1/p}.$$

The constants in both inequalities depend only on d and p.

The proof of this theorem can be found in [6].

Let $p > 0$. Besides the classes $\mathfrak{S}_p$, we will be dealing with the so-called weak Schatten classes Σ_p of compact operators T obeying the condition

$$\|T\|^p_{\Sigma_p} := \sup_{s>0} s^p n(s,T) < \infty.$$

It turns out that Y defined by (2.8) belongs to Σ_p if $a \in L^p$ and the other factor b satisfies the condition

$$\|b\|^p_{L^p_w} := \sup_{s>0}\big(s^p \,\mathrm{measure}\{\xi \in \mathbb{R}^d : \ |b(\xi)| > s\}\big) < \infty.$$

Such functions b are said to belong to the space $L^p_w(\mathbb{R}^d)$. The following result is the so-called Cwikel inequality (see [8]).

Theorem 2.8. *Let $p > 2$. Assume that $a \in L^p(\mathbb{R}^d)$ and $b \in L^p_w(\mathbb{R}^d)$. Then Y defined by* (2.8) *belongs to the class Σ_p and*

$$\|Y\|_{\Sigma_p} \leqslant C\|a\|_{L^p}\|b\|_{L^p_w}$$

with a constant C that depends only on d and p.

3. Preliminary estimates

For the sake of brevity, the norms in the spaces $\mathfrak{S}_p$ and L^p will be often denoted by the symbol $\|\cdot\|_p$.

Theorem 3.1. *Let $l \geqslant 1$ be an integer number. Let p satisfy the condition $p > \frac{6l}{l+1}$. Let also*

$$X_\lambda = W(H_0 - \lambda)^{-1}W$$

be the family of Birman-Schwinger operators with H_0 being the Dirac operator. Assume that $W \in L^p(\mathbb{R}^3)$. Then the operator

$$T_{\lambda,\mu} = X^l_\mu - X^l_\lambda$$

beloings to the Schatten class $\mathfrak{S}_{\frac{p}{2l}}$ and

$$\|T_{\lambda,\mu}\|_{\frac{p}{2l}} \leqslant C\|W\|^{2l}_p, \tag{3.1}$$

with a constant $C > 0$ that does not depend on W.

Proof. It is easy to see that $T_{\lambda,\mu} = X^l_\mu - X^l_\lambda$ is a finite linear combination of operators of the form

$$\Big(WR_\lambda W\Big)^n \Big(WR_\lambda R_\mu W\Big) \Big(WR_\mu W\Big)^m$$

where $R_\lambda = (H_0 - \lambda)^{-1}$ and $n + m = l - 1$. If the factors W were written before the factors R_λ and R_μ, then this term would be the operator

$$W^{2l} R_\lambda^{n+1} R_\mu^{m+1},$$

and Theorem 2.7 would imply an estimate that is similar to (3.1). We have to show that the position of the factors does not matter too much.

For that purpose, we observe that

$$\Big(WR_\lambda W\Big)^n = W|R_\lambda|^{(1+1/l)/2} J_\lambda |R_\lambda|^{(1-1/l)/2} W^{\frac{l-1}{l+1}} W^{\frac{2}{l+1}} \Big(WR_\lambda W\Big)^{n-1}$$

where $J_\lambda = \text{sign}(R_\lambda)$. Consequently,

$$\Big(WR_\lambda W\Big)^n = \Big(\prod_{j=1}^n \Big(W^{q_j}|R_\lambda|^{q_j(1+1/l)/2} J_\lambda |R_\lambda|^{1-q_j(1+1/l)/2} W^{\frac{2l}{l+1}-q_j}\Big)\Big) W^{q_n - \frac{(l-1)}{l+1}}$$

with $q_1 = 1$ and $q_j = q_{j-1} + \frac{2}{l+1} = \frac{l-1+2j}{l+1}$. Therefore,

$$\|\Big(WR_\lambda W\Big)^n W^{-q_n+\frac{(l-1)}{l+1}}\|_r \leqslant \prod_{j=1}^n \|W^{q_j}|R_\lambda|^{q_j(1+1/l)/2}\|_{p/q_j} \||R_\lambda|^{1-q_j(1+1/l)/2} W^{\frac{2l}{l+1}-q_j}\|_{\frac{p(l+1)}{l+1-2j}}$$

with $\frac{1}{r} = \frac{1}{p(l+1)} \sum_{j=1}^n (l - 1 + 2j + (l + 1 - 2j)) = \frac{2ln}{p(l+1)}$. This leads to the estimate

$$\|\Big(WR_\lambda W\Big)^n W^{-q_n+\frac{(l-1)}{l+1}}\|_r \leqslant C \prod_{j=1}^n \|W\|_p^{2l/(l+1)} = C\|W\|_p^{p/r} \tag{3.2}$$

Similarly, since

$$\Big(WR_\mu W\Big)^m = W^{p_1 - \frac{(l-1)}{l+1}} \Big(\prod_{j=1}^m \Big(W^{\frac{2l}{l+1}-p_j}|R_\mu|^{1-p_j(1+1/l)/2} J_\mu |R_\mu|^{p_j(1+1/l)/2} W^{p_j}\Big)\Big)$$

with $p_m = 1$ and $p_{j-1} = p_j + \frac{2}{l+1}$, we obtain that

$$\|W^{-p_1+\frac{(l-1)}{l+1}} \Big(WR_\mu W\Big)^m\|_\tau \leqslant C\|W\|_p^{p/\tau} \tag{3.3}$$

where $\frac{1}{\tau} = \frac{2lm}{p(l+1)}$.

It remains to estimate the Schatten norm of the operator

$$B := W^{q_n+\frac{2}{l+1}} R_\lambda R_\mu W^{p_1+\frac{2}{l+1}} = W^{1+\frac{2n}{l+1}} R_\lambda R_\mu W^{1+\frac{2m}{l+1}}$$

For that purpose, we write it as

$$W^{1+\frac{2n}{l+1}} R_\lambda R_\mu W^{1+\frac{2m}{l+1}} = W^{1+\frac{2n}{l+1}} |R_\lambda|^{(2n+l+1)/(2l)} J_{\lambda,\mu} |R_\mu|^{(2m+l+1)/(2l)} W^{1+\frac{2m}{l+1}}$$

where

$$J_{\lambda,\mu} = |H_0 - \lambda|^{(2n+l+1)/(2l)} R_\lambda R_\mu |H_0 - \mu|^{(2m+l+1)/(2l)}$$

is a bounded operator.

Obviously,

$$\|B\|_\varkappa \leqslant \|J_{\lambda,\mu}\| \, \|W^{1+\frac{2n}{l+1}}|R_\lambda|^{(2n+l+1)/(2l)}\|_{\frac{p(l+1)}{(l+1+2n)}} \||R_\mu|^{(2m+l+1)/(2l)} W^{1+\frac{2m}{l+1}}\|_{\frac{p(l+1)}{(l+1+2m)}}$$

with $\frac{1}{\varkappa} = \frac{2(l+1+n+m)}{p(l+1)} = \frac{2l}{p(l+1)}$. Therefore,

$$\|B\|_{\varkappa} \leqslant C\|W\|_p^{2+\frac{2(n+m)}{(l+1)}} = C\|W\|_p^{p/\varkappa} \tag{3.4}$$

Observe now that

$$\frac{1}{r} + \frac{1}{\tau} + \frac{1}{\varkappa} = \frac{2l}{p}. \tag{3.5}$$

Combining the relations (3.2) -(3.5), we obtain that

$$\|T_{\lambda,\mu}\|_{p/(2l)} \leqslant C\|W\|_p^{2l}.$$

□

If l is an odd number, then

$$T_{\lambda,\mu} = |X_\mu|^l \mathrm{sign}(X_\mu) - |X_\lambda|^l \mathrm{sign}(X_\lambda).$$

In this case, if $p \leqslant 2l$, then it follows from Proposition 2.3 that

$$p\int_0^\infty N(\alpha)\alpha^{-p/2-1}d\alpha \leqslant 2\sum_k s_k^{p/(2l)}(T_{\lambda,\mu}) = 2\|T_{\lambda,\mu}\|_{p/(2l)}^{p/(2l)}.$$

As a consequence, applying Theorem 3.1 with $l = 3$, we obtain Theorem 1.2, saying that

$$q\int_0^\infty N(\alpha)\alpha^{-q-1}d\alpha \leqslant C\|W\|_{2q}^{2q}, \qquad q \in (9/4, 3].$$

4. Splitting

For $\varepsilon > 0$, we introduce two parts V_1 and V_2 of the potential V by setting

$$V_1(x) = \begin{cases} V(x) & \text{if} \quad |x| < \varepsilon \cdot \alpha^{1/\nu}; \\ 0 & \text{if} \quad |x| \geqslant \varepsilon \cdot \alpha^{1/\nu}, \end{cases}$$

and

$$V_2 = V - V_1.$$

Let $N_j(t)$ be the number of eigenvalues of the operator $H_0 - tV_j$ in the interval $[\lambda, \mu)$, $j = 1, 2$. We want to show that

$$\int_0^\alpha N(t)t^{-q-1}dt \sim \int_0^\alpha N_1(t)t^{-q-1}dt + \int_0^\alpha N_2(t)t^{-q-1}dt, \qquad \text{as} \quad \alpha \to \infty.$$

We introduce $\tilde{X}_\lambda$ by

$$\tilde{X}_\lambda = W_1(H_0 - \lambda)^{-1}W_1 + W_2(H_0 - \lambda)^{-1}W_2,$$

where $W_j = \sqrt{V}_j$ for $j = 1, 2$. Note that

$$n_+(t, \tilde{X}_\mu) - n_+(t, \tilde{X}_\lambda) = N_1(t) + N_2(t), \qquad \text{for} \quad 0 < t \leqslant \alpha.$$

As we know from (2.3), for $s = \alpha^{-1}$,

$$\begin{aligned} q\Big|\int_s^\infty \Big(n_+(t, X_\lambda) - n_+(t, \tilde{X}_\lambda)\Big)t^{q-1}dt\Big| &\leqslant \|X_\lambda\|^q + \|\tilde{X}_\lambda\|^q + \\ &+ \sum_{k=1}^{c\alpha^3+1} s_k^{q/3}\Big(X_\lambda^3 - \tilde{X}_\lambda^3\Big). \end{aligned} \tag{4.1}$$

Here, the value of the parameter q is the same as in Theorem 1.2. The next proposition and its corollaries show that the right hand side is of order $o(\alpha^{3/\nu-q})$ as $\alpha\to\infty$. That allows us to replace X_λ and X_μ by the operators $\tilde X_\lambda$ and $\tilde X_\mu$ and claim that

$$\begin{aligned}\int_0^\alpha N(t)t^{-q-1}dt &= \int_{\alpha^{-1}}^\infty \Big(n_+(t,X_\mu)-n_+(t,X_\lambda))\Big)t^{q-1}dt\\ &\sim \int_{\alpha^{-1}}^\infty \Big(n_+(t,\tilde X_\mu)-n_+(t,\tilde X_\lambda))\Big)t^{q-1}dt, \qquad \text{as} \quad \alpha\to\infty.\end{aligned}$$

Proposition 4.1. *Let $p>9/4$ and $\gamma\geqslant 2$. Assume that the support of the function W_2 is contained in the set*

$$\{x\in\mathbb{R}^3:\ \ |x|>\varepsilon\alpha^{1/\nu}+1\}.$$

Let also

$$q=\frac{3\gamma p}{6\gamma+p}.$$

Then there is an $\alpha_0>0$ such that

$$\|X_\lambda^3-\tilde X_\lambda^3\|_{q/3}\leqslant C\|W\|_p^6(\varepsilon^2\alpha^{2/\nu}+1)^{1/\gamma}, \qquad \textit{for} \quad \alpha>\alpha_0,$$

with a constant $C>0$ independent of α and W.

Proof. Let θ be a smooth function on the real line $\mathbb{R}$ such that

$$\theta(t)=\begin{cases}1 & \text{for} \quad t\leqslant 0;\\ 0 & \text{for} \quad t\geqslant 1.\end{cases}$$

Define θ_α on $\mathbb{R}^3$ by

$$\theta_\alpha(x)=\theta(|x|-\varepsilon\alpha^{1/\nu}).$$

Then, obviously, $\theta_\alpha W_1=W_1$ and $\theta_\alpha W_2=0$. Using the identity

$$[B,A^{-1}]=A^{-1}\,[A,B]\,A^{-1},$$

we obtain that

$$\begin{aligned}W_1R_\lambda W_2 &= W_1R_\lambda[H_0,\theta_\alpha]R_\lambda W_2=W_1R_\lambda^2\big[H_0,[H_0,\theta_\alpha]\big]R_\lambda W_2=\\ &= W_1R_\lambda^3\big[H_0,\big[H_0,\big[H_0,\big[H_0,[H_0,\theta_\alpha]\big]\big]\big]\big]R_\lambda^3W_2.\end{aligned}$$

The middle operator $\big[H_0,\big[H_0,\big[H_0,\big[H_0,[H_0,\theta_\alpha]\big]\big]\big]\big]$ is an operator of multiplication by a bounded matrix-valued function supported in the layer

$$\Omega_\alpha=\{x\in\mathbb{R}^3:\ \varepsilon\alpha^{1/\nu}\leqslant|x|\leqslant\varepsilon\alpha^{1/\nu}+1\}.$$

Therefore, the operator

$$Y:=R_\lambda\big[H_0,\big[H_0,\big[H_0,\big[H_0,[H_0,\theta_\alpha]\big]\big]\big]\big]R_\lambda$$

belongs to the Schatten class $\mathfrak{S}_\gamma$ at least for $\gamma\geqslant 2$ and

$$\|Y\|_{\mathfrak{S}_\gamma}^\gamma\leqslant C_0\mathrm{vol}\,\Omega_\alpha\leqslant C(\varepsilon^2\alpha^{2/\nu}+1), \qquad \forall\alpha>0.$$

The operator $X_\lambda^3-\tilde X_\lambda^3$ is a finite linear combination of operators of the form

$$\tilde X_\lambda^n\Big(W_1R_\lambda W_2+W_2R_\lambda W_1\Big)X_\lambda^m$$

where $R_\lambda=(H_0-\lambda)^{-1}$ and $n+m=2$.

Repeating the arguments that lead to the estimate (3.2), we obtain

$$\|\tilde{X}_\lambda^n W^{-\frac{n}{2}}\|_r \leqslant C \prod_{j=1}^{n} \|W\|_p^{3/2} = C\|W\|_p^{p/r} \tag{4.2}$$

with $r = \frac{2p}{3n}$. Similarly, we obtain that

$$\|W^{-\frac{m}{2}} X_\lambda^m\|_\tau \leqslant C\|W\|_p^{p/\tau} \tag{4.3}$$

with $\tau = \frac{2p}{3m}$.

It remains to estimate Schatten norms of the operators

$$B_{1,2} := W_1^{1+\frac{n}{2}} R_\lambda W_2^{1+\frac{m}{2}} \qquad \text{and} \qquad B_{2,1} := W_2^{1+\frac{n}{2}} R_\lambda W_1^{1+\frac{m}{2}}.$$

Obviously, it is enough to estimate only the norm of $B_{1,2}$. For that purpose, we write it as

$$B_{1,2} = W_1^{1+\frac{n}{2}} R_\lambda^2 Y R_\lambda^2 W_2^{1+\frac{m}{2}} = W_1^{1+\frac{n}{2}} |R_\lambda|^{(n+2)/3} Q_\lambda |R_\mu|^{(m+2)/3} W_2^{1+\frac{m}{2}}$$

where

$$Q_\lambda = |H_0 - \lambda|^{(n+2)/3} R_\lambda^2 Y R_\lambda^2 |H_0 - \lambda|^{(m+2)/3}$$

belongs to $\mathfrak{S}_\gamma$ and $\|Q_\lambda\|_{\mathfrak{S}_\gamma} \leqslant C\|Y\|_{\mathfrak{S}_\gamma}$.

Obviously,

$$\|B_{1,2}\|_\varkappa \leqslant \|W^{1+\frac{n}{2}} |R_\lambda|^{(n+2)/3}\|_{\frac{2p}{(2+n)}} \||R_\lambda|^{(m+2)/3} W^{1+\frac{m}{2}}\|_{\frac{2p}{(2+m)}} \|Q_\lambda\|_\gamma$$

with $\frac{1}{\varkappa} = \frac{3}{p} + \frac{1}{\gamma}$. Therefore,

$$\|B_{1,2}\|_\varkappa \leqslant C\|W\|_p^3 (\varepsilon^2 \alpha^{2/\nu} + 1)^{1/\gamma} \tag{4.4}$$

Observe now that

$$\frac{1}{r} + \frac{1}{\tau} + \frac{1}{\varkappa} = \frac{6}{p} + \frac{1}{\gamma} = \frac{3}{q}. \tag{4.5}$$

Combining the relations (4.2) -(4.5), we obtain that

$$\|X_\lambda^3 - \tilde{X}_\lambda^3\|_{q/3} \leqslant C\|W\|_p^6 (\varepsilon^2 \alpha^{2/\nu} + 1)^{1/\gamma}.$$

□

In fact we proved more: exactly the same arguments can be used to justify the following statement.

Corollary 4.2. *Let $p > 9/4$, $\gamma \geqslant 2$ and*

$$q = \frac{3\gamma p}{6\gamma + p}.$$

Let the operator $T(\alpha)$ be a finite linear combination of products of three factors of the form

$$\chi_j^- X_\lambda \chi_j^+, \qquad j = 1, 2, 3, \tag{4.6}$$

where $\chi_j^\pm$ are characteristic functions of some subsets of $\mathbb{R}^3$ (that might depend on α). Assume that, at least for one of the three factors (4.6) *in each product, the supports of χ_j^- and χ_j^+ are separated from each other by a spherical layer of the form*

$$\{x \in \mathbb{R}^3 : \ \varepsilon\alpha^{1/\nu} + a \leqslant |x| \leqslant \varepsilon\alpha^{1/\nu} + b\}, \qquad \textit{with} \quad a < b.$$

Then there is an $\alpha_0 > 0$ such that

$$\|T(\alpha)\|_{q/3} \leqslant C\|W\|_p^6 (\varepsilon^2 \alpha^{2/\nu} + 1)^{1/\gamma}, \qquad \textit{for} \quad \alpha > \alpha_0,$$

with a constant $C > 0$ independent of α and W.

As a consequence, we immediately obtain the next result, in which we use the notation $R_\lambda = (H_0 - \lambda)^{-1}$.

Corollary 4.3. *Let $p > 9/4$, $\gamma \geqslant 2$ and*

$$q = \frac{3\gamma p}{6\gamma + p}.$$

Let χ_j be the characteristic functions of the sets

$$\{x \in \mathbb{R}^3 : \ \varepsilon\alpha^{1/\nu} - j \leqslant |x| \leqslant \varepsilon\alpha^{1/\nu} + j\}, \qquad j = 1, 2, 3,$$

and $Y_{\lambda,\alpha}$ be the operator defined by

$$\begin{aligned} Y_{\lambda,\alpha} = \chi_3 \tilde{X}_\lambda \chi_2 \tilde{X}_\lambda \chi_1 \Big(W_1 R_\lambda W_2 + W_2 R_\lambda W_1 \Big) \chi_1 + \chi_2 \tilde{X}_\lambda \chi_1 \Big(W_1 R_\lambda W_2 + \\ W_2 R_\lambda W_1 \Big) \chi_1 X_\lambda \chi_2 + \chi_1 \Big(W_1 R_\lambda W_2 + W_2 R_\lambda W_1 \Big) \chi_1 X_\lambda \chi_2 X_\lambda \chi_3. \end{aligned} \tag{4.7}$$

Then there is an $\alpha_0 > 0$ such that

$$\|X_\lambda^3 - \tilde{X}_\lambda^3 - Y_{\lambda,\alpha}\|_{q/3} \leqslant C\|W\|_p^6 (\varepsilon^2 \alpha^{2/\nu} + 1)^{1/\gamma}, \qquad \text{for} \quad \alpha > \alpha_0,$$

with a constant $C > 0$ independent of α and W.

Proof. One only needs to realize that the operator $T(\alpha) := X_\lambda^3 - \tilde{X}_\lambda^3 - Y_{\lambda,\alpha}$ satisfies conditions of Corollary 4.2. □

On the other hand, applying Cwikel's inequality, one can easily show that

$$\|Y_{\lambda,\alpha}\|_{\Sigma_1} \leqslant C \int_{\mathbb{R}^3} \chi_3(x) V^3(x) dx \leqslant C_\varepsilon \alpha^{2/\nu - 3} \qquad \text{for} \quad \alpha > \alpha_0.$$

In other words,

$$s_k\big(Y_{\lambda,\alpha}\big) \leqslant C_\varepsilon \alpha^{2/\nu-3} k^{-1} \qquad \text{for} \quad \alpha > \alpha_0.$$

Consequently,

$$\sum_1^{c\alpha^3+1} s_k^{q/3}\big(Y_{\lambda,\alpha}\big) \leqslant C_\varepsilon \alpha^{2/\nu - q}. \tag{4.8}$$

Corollary 4.4. *Let $\frac{9}{4} < q < \frac{3}{\nu}$ where $\nu > 1$. Then*

$$\int_{\alpha^{-1}}^{\infty} \Big(n_+(t, X_\lambda) - n_+(t, \tilde{X}_\lambda) \Big) t^{q-1} dt = o(\alpha^{3/\nu - q}), \qquad \text{as} \quad \alpha \to \infty.$$

Proof. Choose $\gamma > \frac{q}{3-\nu q}$ and define p by

$$\frac{6}{p} = \frac{3}{q} - \frac{1}{\gamma}.$$

Then $p > 6/\nu$. Therefore $W \in L^p(\mathbb{R}^3)$. Using Rotfeld's inequality, we obtain

$$\sum_1^{c\alpha^3+1} s_k^{q/3}\big(X_\lambda^3 - \tilde{X}_\lambda^3\big) \leqslant \sum_1^{c\alpha^3+1} s_k^{q/3}\big(Y_{\lambda,\alpha}\big) + \|X_\lambda^3 - \tilde{X}_\lambda^3 - Y_{\lambda,\alpha}\|_{q/3}^{q/3}. \tag{4.9}$$

The inequality (4.9) estimates the last term on the right hand side of (4.1). It remains to apply Corollary 4.3 and the relation (4.8). □

5. Other consequences

The preceding discussion of the splitting principle, involves a decomposition of the space $\mathbb{R}^3$ into two domains. It is easy to see that the same arguments work for all piecewise smooth domains obtained similarly by scaling by a factor of $\alpha^{1/\nu}$. In particular, one of the domains that we have alredy considered can be decomposed further into smaller sets. Namely, let Q_j be bounded disjoint cubes contained in the region $\{x \in \mathbb{R}^3 : \ |x| \geqslant \varepsilon\}$, $1 \leqslant j \leqslant n-1$. Let $\{\phi_j\}_{j=1}^{n-1}$ be the characteristic functions of the cubes $\alpha^{1/\nu}Q_j$. Define ϕ_n to be the characteristic function of the complement

$$\{x \in \mathbb{R}^3 : \ |x| \geqslant \varepsilon\alpha^{1/\nu}\} \setminus \cup_{j=1}^{n-1}\alpha^{1/\nu}Q_j.$$

Theorem 5.1. *Let $\frac{9}{4} < q < \frac{3}{\nu}$ where $\nu > 1$. Then*

$$\int_{\alpha^{-1}}^{\infty}\Big(n_+\big(t, W_2(H_0-\lambda)^{-1}W_2\big) - \sum_{j=1}^{n} n_+\big(t, \phi_j W(H_0-\lambda)^{-1}W\phi_j\big)\Big)t^{q-1}dt = o(\alpha^{3/\nu-q}),$$

as $\alpha \to \infty$.

To prove Theorem 5.1, it is enough to repeat the steps that were needed to prove Corollary 4.4.

Clearly, to obtain an asymptotic formula for $\int_{\alpha^{-1}}^{\infty} n_+\big(t, W_2(H_0-\lambda)^{-1}W_2\big)t^{q-1}dt$ one has to obtain an asymtotic formula for $\int_{\alpha^{-1}}^{\infty} n_+\big(t, \phi_j W(H_0-\lambda)^{-1}W\phi_j\big)t^{q-1}dt$ for each j. The latter integral can be written as

$$\int_{\alpha^{-1}}^{\infty} n_+\big(t, \phi_j W(H_0-\lambda)^{-1}W\phi_j\big)t^{q-1}dt = \alpha^{-q}\int_{1}^{\infty} n_+\Big(\frac{\tau}{\alpha}, \phi_j W(H_0-\lambda)^{-1}W\phi_j\Big)\tau^{q-1}d\tau$$

Observe now that, if

$$V(x) = \frac{\Phi(\theta)}{|x|^\nu}, \qquad \text{for} \quad |x| > 1,$$

then the maximum and minimum values of αV on the cubes $\alpha^{1/\nu}Q_j$ do not depend on α:

$$m_j \leqslant \alpha V(x) \leqslant M_j, \qquad \text{for all} \quad \alpha > \varepsilon^{-\nu} \quad \text{and all} \quad x \in \alpha^{1/\nu}Q_j.$$

The potential αV can be estimated by constant functions $m_j\phi_j$ and $M_j\phi_j$ on cubes $\alpha^{1/\nu}Q_j$.

So, due to the monotonicity of the counting function n_+,

$$n_+\left(\frac{\tau}{m_j}, \phi_j(H_0-\lambda)^{-1}\phi_j\right) \leqslant n_+\left(\frac{\tau}{\alpha}, \phi_j W(H_0-\lambda)^{-1}W\phi_j\right) \leqslant n_+\left(\frac{\tau}{M_j}, \phi_j(H_0-\lambda)^{-1}\phi_j\right),$$

for any $\tau > 0$.

Consequently, it remains to obtain an asymptotic formula for the quantity

$$n_+\big(t, \phi_j(H_0-\lambda)^{-1}\phi_j\big) \quad \text{as} \quad \alpha \to \infty$$

for any fixed $t > 0$. We are going to prove the following result.

Proposition 5.2. *For any fixed $t > 0$,*

$$n_+\big(t, \phi_j(H_0-\lambda)^{-1}\phi_j\big) \sim 3^{-1}\pi^{-2}\alpha^{3/\nu}\Big((t^{-1}+\lambda)_+^2 - 1\Big)_+^{3/2}\operatorname{vol} Q_j \quad \textit{as} \quad \alpha \to \infty.$$

Proof. Note that

$$\frac{4}{3}\pi\Big((t^{-1}+\lambda)_+^2-1\Big)_+^{3/2}=\mathrm{vol}\,\{\xi\in\mathbb{R}^3:\quad \big(\sqrt{|\xi|^2+1}-\lambda\big)^{-1}>t\}.$$

Taking into account the fact that $\pm\sqrt{|\xi|^2+1}$ are eigenvalues of the symbol

$$A(\xi):=\sum_1^3\gamma_j\xi_j+\gamma_0$$

of the operator H_0, we conclude that we need to prove that

$$\mathrm{tr}\,\Psi\big(\phi_j(H_0-\lambda)^{-1}\phi_j\big)\sim\mathrm{tr}\,\Big[\phi_j\Psi\big((H_0-\lambda)^{-1}\big)\phi_j\Big]\quad\text{as}\quad\alpha\to\infty,\tag{5.1}$$

for Ψ being the characteristic function of the interval (t,∞). Since such a function Ψ can be estimated from above and below by continuous functions of the form

$$\Psi_\epsilon(s)=\begin{cases}0, & \text{if}\quad s<\tau\\ (s-\tau)/\epsilon, & \text{if}\quad \tau\leqslant s\leqslant\tau+\epsilon\\ 1, & \text{if}\quad s>\tau+\epsilon,\end{cases}$$

and the quantity

$$\Big((t^{-1}+\lambda)_+^2-1\Big)_+^{3/2}$$

depends on t continuously, we only need to prove (5.1) for Ψ that are continuous and vanishing near zero.

Any such function Ψ can be written as

$$\Psi(s)=s^5\zeta(s),$$

where ζ is a continuous function on the real line $\mathbb{R}$. Notice that, in this case,

$$\Big|\mathrm{tr}\,\Psi\big(\phi_j(H_0-\lambda)^{-1}\phi_j\big)\Big|\leqslant\|\phi_j(H_0-\lambda)^{-1}\phi_j\|_{\mathfrak{S}_5}^5\|\zeta\|_\infty\leqslant C\alpha^{3/\nu}\|\zeta\|_\infty.$$

Moreover,

$$\Big|\mathrm{tr}\,\phi_j\Psi\big((H_0-\lambda)^{-1}\big)\phi_j\Big|\leqslant\|\phi_j(H_0-\lambda)^{-2}\|_{\mathfrak{S}_{5/2}}\|(H_0-\lambda)^{-3}\phi_j\|_{\mathfrak{S}_{5/3}}\|\zeta\|_\infty\leqslant C\alpha^{3/\nu}\|\zeta\|_\infty,$$

Thus both sides of (5.1) can be estimated by $C\alpha^{3/\nu}\|\zeta\|_\infty$. The functional $\|\zeta\|_\infty$ is the L^∞-norm of the function on the interval $[-L,L]$ where $L=1/(1-|\lambda|)$. Since ζ can be uniformly approximated by polynomials, it is enough to prove (5.1) under the assumption that ζ is a polynomial. Put differently, it is enough to prove it for

$$\Psi(s)=s^n,\qquad n\geqslant 5.$$

Denote $B=(H_0-\lambda)^{-1}$, $\chi_+=\phi_j$ and $\chi_-=1-\phi_j$. We are going to prove that

$$\|(\chi_+B\chi_+)^n-\chi_+B^n\chi_+\|_{\mathfrak{S}_1}=o(\alpha^{3/\nu}),\qquad\text{as}\quad\alpha\to\infty.\tag{5.2}$$

For that purpose, we write $\chi_+B^n\chi_+$ as

$$\chi_+B^n\chi_+=(\chi_+B\chi_+)^n+\sum_{j=0}^{n-1}(\chi_+B\chi_+)^j\chi_+B\chi_-B^{n-j-1}\chi_+.\tag{5.3}$$

While the norm of the operator $\chi_+B\chi_-$ does not tend to zero, it is still representable in the form

$$\chi_+B\chi_-=T_1+T_2,\qquad\text{where}\qquad \|T_1\|\to 0,\quad\text{and}\quad \|T_2\|_{\mathfrak{S}_n}=o(\alpha^{3/(n\nu)}),\quad\text{as}\quad\alpha\to\infty.$$

To see that, we define T_2 to be the operator

$$T_2 = \theta\chi_+ B\chi_-\theta,$$

where θ is the operator of multiplication by the characteristic function of the set

$$\big((\alpha^{1/\nu} + \alpha^{1/(2\nu)})Q_j\big) \setminus \big((\alpha^{1/\nu} - \alpha^{1/(2\nu)})Q_j\big).$$

Then the volume of the support of the function θ does not exceed $C\alpha^{5/(2\nu)}$. Therefore,

$$\|T_2\|_{\mathfrak{S}_n} \leqslant C\alpha^{5/(2n\nu)} = o(\alpha^{3/(n\nu)}), \quad \text{as} \quad \alpha \to \infty.$$

On the other hand, we have the estimate for the integrtal kernel $k(x,y)$ of the operator T_1:

$$|k(x,y)| \leqslant C(1-\theta(x))(1-\theta(y))e^{-c|x-y|}$$

which implies that $\|T_1\| \to 0$ as $\alpha \to \infty$ because x and y are getting far away from each other while

$$\|T_1\| \leqslant \Big(\sup_x \int |k(x,y)|dy \times \sup_y \int |k(x,y)|dx\Big)^{1/2}.$$

Thus, we have the following estimate

$$\begin{aligned} \|(\chi_+B\chi_+)^j\chi_+B\chi_-B^{n-j-1}\chi_+\|_{\mathfrak{S}_1} \leqslant \|(\chi_+B\chi_+)^jT_1B^{n-j-1}\chi_+\|_{\mathfrak{S}_1} + \\ \|(\chi_+B\chi_+)^jT_2B^{n-j-1}\chi_+\|_{\mathfrak{S}_1} \leqslant \|\chi_+B\chi_+\|^j_{\mathfrak{S}_{n-1}}\|T_1\|\|B^{n-j-1}\chi_+\|_{\mathfrak{S}_{(n-1)/(n-j-1)}} \\ +\|\chi_+B\chi_+\|^j_{\mathfrak{S}_n}\|T_2\|_{\mathfrak{S}_n}\|B^{n-j-1}\chi_+\|_{\mathfrak{S}_{n/(n-j-1)}} = o(\alpha^{3/\nu}), \qquad \text{as} \qquad \alpha \to \infty. \end{aligned}$$

Combining this relation with (5.3) we obtain (5.2). $\square$

As a consequence, we obtain

Proposition 5.3. *For any constant $M \geqslant 0$, we have*

$$\begin{aligned} \lim_{\alpha\to\infty} \alpha^{-3/\nu+q} \int_{\alpha^{-1}}^{\infty} n_+\big(\alpha t, M\phi_j(H_0-\lambda)^{-1}\phi_j\big)t^{q-1}dt = \\ 3^{-1}\pi^{-2}\mathrm{vol}\, Q_j \int_1^\infty \Big((t^{-1}M+\lambda)_+^2 - 1\Big)_+^{3/2} t^{q-1}dt \end{aligned} \tag{5.4}$$

Proof. Changing the variables $\alpha t \to t$, we obtain

$$\begin{aligned} \lim_{\alpha\to\infty} \alpha^{-3/\nu+q} \int_{\alpha^{-1}}^{\infty} n_+\big(\alpha t, M\phi_j(H_0-\lambda)^{-1}\phi_j\big)t^{q-1}dt = \\ \lim_{\alpha\to\infty} \alpha^{-3/\nu} \int_1^\infty n_+\big(t, M\phi_j(H_0-\lambda)^{-1}\phi_j\big)t^{q-1}dt \end{aligned} \tag{5.5}$$

The integrand on the right hand side can be estimated according to Cwikel's inequality:

$$n_+(t, M\phi_j(H_0-\lambda)^{-1}\phi_j) \leqslant Ct^{-3}\int_{\mathbb{R}^3} \phi_j^6(x)dx \leqslant C\alpha^{3/\nu}t^{-3}\mathrm{vol}\, Q_j.$$

Consequently, the limit on the right hand side of (5.5) can be computed by the Lebesgue dominated convergence theorem. The relation (5.4) follows now from Proposition 5.2. $\square$

To state the next result, we need to introduce the notation

$$\mathfrak{T}(M,\lambda) = 3^{-1}\pi^{-2}\int_1^\infty \Big((t^{-1}M+\lambda)_+^2 - 1\Big)_+^{3/2} t^{q-1}dt$$

Theorem 5.4. *Assume that*

$$V(x) = \frac{\Phi(\theta)}{|x|^\nu}, \qquad \textit{for} \qquad |x| > 1, \tag{5.6}$$

where Φ is a continuous function on the unit sphere. Then

$$\lim_{\alpha\to\infty} \alpha^{-3/\nu+q} \int_{\alpha^{-1}}^\infty n_+\big(t, W_2(H_0-\lambda)^{-1}W_2\big) t^{q-1} dt = \int_{|x|>\varepsilon} \mathfrak{T}\big(\Phi(\theta)|x|^{-\nu}, \lambda\big)\, dx. \tag{5.7}$$

Proof. Let m_j and M_j be the maximum and the minimum values of V on the cube Q_j. Then, according to Proposition 5.3, we have

$$\begin{gathered} 3^{-1}\pi^{-2} \mathrm{vol}\, Q_j \int_1^\infty \Big((t^{-1}m_j+\lambda)_+^2 - 1\Big)_+^{3/2} t^{q-1} dt \leqslant \\ \lim_{\alpha\to\infty} \alpha^{-3/\nu+q} \int_{\alpha^{-1}}^\infty n_+\big(t, \phi_j W(H_0-\lambda)^{-1} W\phi_j\big) t^{q-1} dt \leqslant \\ 3^{-1}\pi^{-2} \mathrm{vol}\, Q_j \int_1^\infty \Big((t^{-1}M_j+\lambda)_+^2 - 1\Big)_+^{3/2} t^{q-1} dt, \end{gathered} \tag{5.8}$$

by the monotonicity of the counting function n_+. Taking the sum over j on the three sides of (5.8) and using Theorem 5.1, we obtain that

$$\sum_{j=1}^n \mathfrak{T}(m_j,\lambda) \mathrm{vol}\, Q_j \leqslant \lim_{\alpha\to\infty} \alpha^{-3/\nu+q} \int_{\alpha^{-1}}^\infty n_+\big(t, W_2(H_0-\lambda)^{-1}W_2\big) t^{q-1} dt \leqslant \sum_{j=1}^n \mathfrak{T}(M_j,\lambda) \mathrm{vol}\, Q_j,$$

It remains to realize that the left and the right hand sides are the Riemann sums of the integral on the right hand side of (5.7). $\square$

Obviously, the condition (5.6) of the last theorem can be replaced by the assumption that the right hands side is only the asymptotics of V.

Theorem 5.5. *Let $V \geqslant 0$ be a bounded real valued potential such that*

$$V(x) = \frac{\Phi(\theta)}{|x|^\nu}\Big(1+o(1)\Big), \qquad \textit{as} \qquad |x|\to\infty,$$

where Φ is a continuous function on the unit sphere. Let $9/4 < q < 3/\nu$ and $\nu > 1$. Then

$$\begin{gathered} \lim_{\alpha\to\infty} \alpha^{-3/\nu+q} \int_{\alpha^{-1}}^\infty n_+(t, W_2(H_0-\lambda)^{-1}W_2) t^{q-1} dt = \\ = 3^{-1}\pi^{-2} \int_1^\infty \Big(\int_{|x|>\varepsilon} \Big((t^{-1}\Phi(\theta)|x|^{-\nu}+\lambda)_+^2 - 1\Big)_+^{3/2} dx\Big) t^{q-1} dt. \end{gathered} \tag{5.9}$$

6. The end of the proof

Proposition 6.1. *Let $V \geqslant 0$ be a bounded real valued potential such that*

$$V(x) = \frac{\Phi(\theta)}{|x|^\nu}\Big(1 + o(1)\Big), \qquad as \qquad |x| \to \infty,$$

where Φ is a continuous function on the unit sphere. Let $9/4 < q < 3/\nu$ and $\nu > 1$. Let also $-1 < \lambda < \mu < 1$. Then

$$\limsup_{\alpha\to\infty} \alpha^{-3/\nu+q} \int_{\alpha^{-1}}^\infty \Big(n_+(t, W_1(H_0-\mu)^{-1}W_1) - n_+(t, W_1(H_0-\lambda)^{-1}W_1)\Big)t^{q-1}dt$$
$$\leqslant \quad \frac{4\pi\varepsilon^{3-\nu q}}{3-\nu q}\|\Phi\|_\infty. \tag{6.1}$$

Proof. It is enough to apply the estimate established in Theorem 1.2 with V replaced by the potential V_1. $\square$

Corollary 6.2. *Let $V \geqslant 0$ be a bounded real valued potential such that*

$$V(x) = \frac{\Phi(\theta)}{|x|^\nu}\Big(1 + o(1)\Big), \qquad as \qquad |x| \to \infty,$$

where Φ is a continuous function on the unit sphere. Let $9/4 < q < 3/\nu$ and $\nu > 1$. Let also $-1 < \lambda < \mu < 1$. Then

$$3^{-1}\pi^{-2}\int_1^\infty \Big(\int_{|x|>\varepsilon} \big(((t^{-1}\Phi(\theta)|x|^{-\nu} + \mu)_+^2 - 1)_+^{3/2}$$
$$-\big((t^{-1}\Phi(\theta)|x|^{-\nu} + \lambda)_+^2 - 1\big)_+^{3/2}\big)dx\Big)t^{q-1}dt \leqslant \tag{6.2}$$
$$\liminf_{\alpha\to\infty} \alpha^{-3/\nu+q}\int_{\alpha^{-1}}^\infty \Big(n_+(t, \tilde X_\mu) - n_+(t, \tilde X_\lambda)\Big)t^{q-1}dt,$$

while

$$\limsup_{\alpha\to\infty} \alpha^{-3/\nu+q}\int_{\alpha^{-1}}^\infty \Big(n_+(t, \tilde X_\mu) - n_+(t, \tilde X_\lambda)\Big)t^{q-1}dt \leqslant \quad \frac{4\pi\varepsilon^{3-\nu q}}{3-\nu q}\|\Phi\|_\infty +$$
$$3^{-1}\pi^{-2}\int_1^\infty \Big(\int_{|x|>\varepsilon} \big(((t^{-1}\Phi(\theta)|x|^{-\nu} + \mu)_+^2 - 1)_+^{3/2} \tag{6.3}$$
$$-\big((t^{-1}\Phi(\theta)|x|^{-\nu} + \lambda)_+^2 - 1\big)_+^{3/2}\big)dx\Big)t^{q-1}dt.$$

Theorem 6.3. *Let $V \geqslant 0$ be a bounded real valued potential such that*

$$V(x) = \frac{\Phi(\theta)}{|x|^\nu}\Big(1 + o(1)\Big), \qquad as \qquad |x| \to \infty,$$

where Φ is a continuous function on the unit sphere. Let $9/4 < q < 3/\nu$ and $\nu > 1$. Let also $-1 < \lambda < \mu < 1$. Then

$$\lim_{\alpha\to\infty} \alpha^{-3/\nu+q}\int_{\alpha^{-1}}^\infty \Big(n_+(t, X_\mu) - n_+(t, X_\lambda)\Big)t^{q-1}dt$$
$$= 3^{-1}\pi^{-2}\int_1^\infty \Big(\int_{\mathbb{R}^3}\Big[\Big((t^{-1}\Phi(\theta)|x|^{-\nu} + \mu)_+^2 - 1\Big)_+^{3/2} - \tag{6.4}$$
$$\Big((t^{-1}\Phi(\theta)|x|^{-\nu} + \lambda)_+^2 - 1\Big)_+^{3/2}\Big]dx\Big)t^{q-1}dt.$$

Proof. According to Corollary 4.4, $\tilde{X}_\lambda$ and $\tilde{X}_\mu$ in (6.2) and (6.3) can be replaced by the operatorsX_λ and X_μ. After this replacement, we can pass to the limit as $\varepsilon \to 0$. □

Theorem 1.1 is now a consequence of Theorem 6.3.

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